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**INITIAL BOUNDARY VALUE PROBLEM FOR FRACTAL  
HEAT EQUATION IN THE SEMI-INFINITE REGION  
BY YANG-LAPLACE TRANSFORM**

by

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Short paper

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*Analytical solution of transient heat conduction through a semi-infinite fractal medium is developed. The solution focuses on application of a local fractional derivative operator to model the heat transfer process and a solution through the Yang-Laplace transform.*

Key words: *initial boundary value problem, heat equation, fractal media, Yang-Laplace transform*

## Introduction

Boundary value problems in heat conduction [1-4] are ever attractive due to their significant practical and academic aspects and commonly solved by numerical [5] or analytical [6] techniques such as variational iteration methods [7], Adomian's decomposition method [8] and the homotopy analysis method [9].

In case of fractional difference models of heat conduction [10-14] describing anomalous transport of thermal energy [15] in fractal media, the boundary value problems are described by fractional diffusion equations [16, 17] solved numerically or analytically [18-21]. Moreover, especially in the case of heat conduction, this leads to non-differentiable transport problems [22-27] solved in a variety ways, among them: local fractional variation iteration method [22, 23], local fractional Fourier series method [26], local fractional Laplace variational iteration method [27], etc.

This communication addresses a solution of transient heat conduction problem through a semi-infinite fractal medium [26] and developed by the Yang-Laplace transform. The Yang-Laplace transform [27-29] was conceived to solve some differential equations expressed through local fractional derivatives.

## The mathematical method

For seek of clarity of the explanation, the properties for Yang-Laplace transform are briefly outlined. The Yang-Laplace transform is defined as [27-29]:

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$$\tilde{L}_\alpha \{f(x)\} = f_s^{\tilde{L},\alpha}(s) = \frac{1}{\Gamma(1+\alpha)} \int_0^\infty E_\alpha(-s^\alpha x^\alpha) f(x) (dx)^\alpha, \quad 0 < \alpha \leq 1 \quad (1)$$

and its inverse formula is defined as [27-29]:

$$f(x) = \tilde{L}_\alpha^{-1} \{f_s^{\tilde{L},\alpha}(s)\} = \frac{1}{(2\pi)^\alpha} \int_{\beta-i\infty}^{\beta+i\infty} E_\alpha(s^\alpha x^\alpha) f_s^{\tilde{L},\alpha}(s) (ds)^\alpha \quad (2)$$

where  $f(x)$  is the local fractional continuous,  $s^\alpha = \beta^\alpha + i^\alpha \infty^\alpha$  and  $\text{Re}(s^\alpha) = \beta^\alpha$ .

The following properties for Yang-Laplace transform are valid [28]:

$$\begin{aligned} \tilde{L}_\alpha \{f^{(\omega)}(x)\} &= s^\alpha \tilde{L}_\alpha \{f(x)\} - f(0), \quad \tilde{L}_\alpha \{x^{k\alpha}\} = \frac{\Gamma(1+k\alpha)}{s^{(k+1)\alpha}}, \\ L_\alpha \{f(ax)\} &= \frac{1}{a^\alpha} f_s^{\tilde{L},\alpha}\left(\frac{s}{a}\right), \quad a > 0 \end{aligned} \quad (3a,b,c)$$

For more details of local fractional derivatives and integrals, see [26, 28].

### Fractal heat equation in the semi-infinite region and its solution

The first law of thermodynamic states in fractal media reads as [26]:

$$\frac{1}{\Gamma(1+\alpha)} \int_{\tau_0}^{\tau} \left[ \int_{V(\tau)} \left( K^{2\alpha} \nabla^\alpha u + g - \rho_\alpha c_\alpha \frac{\partial^\alpha u}{\partial t^\alpha} \right) dV(\tau) \right] (dt)^\alpha = 0 \quad (4a)$$

which leads to the equation [25, 26]:

$$K^{2\alpha} \nabla^\alpha u + g - \rho_\alpha c_\alpha \frac{\partial^\alpha u}{\partial t^\alpha} = 0 \quad (4b)$$

where the volume integral is the local fractional volume integral [26]. When the fractal dimension is equal to 1, eq. (4b) becomes the classical Fourier equation.

Making use of  $g = 0$  and  $\rho_\alpha c_\alpha = K^{2\alpha}$ , the 1-D heat conduction through a semi-infinite fractal medium is modelled by [26, 27, 30]:

$$\frac{\partial^\alpha u(x,t)}{\partial t^\alpha} - \frac{\partial^{2\alpha} u(x,t)}{\partial x^{2\alpha}} = 0, \quad x > 0, \quad t > 0 \quad (5a)$$

$$u(x,0) = 0, \quad u(0,t) = u_0 \quad (5b,c)$$

with the Yang-Laplace transform, the model (5a, b, c) can be transformed into:

$$s^\alpha u(x,s) - u(x,0) - \frac{\partial^{2\alpha} u(x,s)}{\partial x^{2\alpha}} = 0, \quad u(0,s) = \frac{u_0}{s^\alpha} \quad (6a,b)$$

From eq. (6a), we get:

$$s^\alpha u(x,s) - \frac{\partial^{2\alpha} u(x,s)}{\partial x^{2\alpha}} = 0 \quad (7)$$

where the initial value condition is eq. (6b).

The general solution of eq. (7) can be expressed in the form:

$$u(x,s) = A E_\alpha(s^{\frac{\alpha}{2}} x^\alpha) + B E_\alpha(-s^{\frac{\alpha}{2}} x^\alpha) \quad (8)$$

In the expression (8) the pre-factors  $A$  and  $B$  are constants. However, taking into account that the temperature function is bounded, we get  $A = 0$ . Hence, from eq. (8) we have:

$$u(x, s) = \frac{u_0}{s^\alpha} E_\alpha(-s^{\frac{\alpha}{2}} x^\alpha) \quad (9)$$

Taking into account the transforms:

$$\tilde{L}_\alpha \{t^{2\alpha} f(t)\} = \frac{\partial^{2\alpha} T(s)}{\partial s^{2\alpha}}, \quad \tilde{L}_\alpha \{-t^\alpha f(t)\} = \frac{\partial^\alpha T(s)}{\partial s^\alpha}, \quad (10 \text{ a,b,c})$$

$$\tilde{L}_\alpha \{f(t)\} = T(s) = E_\alpha(-s^{\frac{\alpha}{2}})$$

we may obtain:

$$\frac{\partial^\alpha f(t)}{\partial t^\alpha} + mf(t) = 0, \quad m = \frac{4^\alpha t^\alpha \mu - 1}{4^\alpha t^{2\alpha}}, \quad \mu = \frac{\Gamma(1+2\alpha)}{\Gamma(1+\alpha)} - \frac{\Gamma\left(1-\frac{\alpha}{2}\right)}{\Gamma\left(1-\frac{3\alpha}{2}\right)} \quad (10 \text{ d,e,f})$$

Following (10d) we have:

$$f(t) = \frac{1}{2^\alpha \Gamma\left(1-\frac{\alpha}{2}\right) t^{\frac{3\alpha}{2}}} E_\alpha \left[ -\frac{\Gamma\left(1-\frac{5\alpha}{2}\right)}{4^\alpha \mu \Gamma\left(1-\frac{3\alpha}{2}\right) t^\alpha} \right] \quad (11)$$

Applying the Yang-Laplace transform we have:

$$\tilde{L}_\alpha^{-1} \left\{ E_\alpha \left( -s^{\frac{\alpha}{2}} \right) \right\} = \frac{1}{2^\alpha \Gamma\left(1-\frac{\alpha}{2}\right) t^{\frac{3\alpha}{2}}} E_\alpha \left[ -\frac{\Gamma\left(1-\frac{5\alpha}{2}\right)}{4^\alpha \mu \Gamma\left(1-\frac{3\alpha}{2}\right) t^\alpha} \right] \quad (12a)$$

and

$$\tilde{L}_\alpha^{-1} \left\{ \frac{E_\alpha \left( -s^{\frac{\alpha}{2}} x^\alpha \right)}{x^{2\alpha} s^\alpha} \right\} = \frac{1}{2^\alpha \Gamma\left(1-\frac{\alpha}{2}\right) \Gamma(1+\alpha)} \int_0^t \frac{1}{\left(\frac{\tau}{x^2}\right)^{\frac{3\alpha}{2}}} E_\alpha \left[ -\frac{\Gamma\left(1-\frac{5\alpha}{2}\right)}{4^\alpha \mu \Gamma\left(1-\frac{3\alpha}{2}\right) \left(\frac{\tau}{x^2}\right)^\alpha} \right] (d\tau)^\alpha \quad (12b)$$

From eq. (12b) we obtain the non-differentiable solution of eq. (5a) in the form:

$$u(x, t) = \frac{u_0 x^{2\alpha}}{2^\alpha \Gamma\left(1-\frac{\alpha}{2}\right) \Gamma(1+\alpha)} \cdot$$

$$\int_0^t \frac{1}{\left(\frac{\tau}{x^2}\right)^{\frac{3\alpha}{2}}} E_\alpha \left\{ \frac{\Gamma\left(1-\frac{5\alpha}{2}\right)}{4^\alpha \left(\frac{\tau}{x^2}\right)^\alpha \left[ \frac{\Gamma(1+2\alpha)}{\Gamma(1+\alpha)} - \frac{\Gamma\left(1-\frac{\alpha}{2}\right)}{\Gamma\left(1-\frac{3\alpha}{2}\right)} \right] \Gamma\left(1-\frac{3\alpha}{2}\right)} \right\} (d\tau)^\alpha \quad (13)$$

The solution (13) is a fractal function in accordance with the local fractional continuity concept [26].

## Conclusions

The Yang-Laplace transform was successfully applied to solve an initial boundary value problem for fractal heat equation in the semi-infinite region, with local fractional derivatives and non-differentiable conditions. The result differs from those developed in [10-14] due to differences in the employed fractal operators. The solution allows, when the fractal dimension is  $\alpha = \ln 2/\ln 3$  a solution on Cantor sets to be developed.

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## Nomenclature

$c_\alpha$  – specific heat of fractal material, [ $\text{Jkg}^{-1}$ ]  
 $g$  – energy generation term  
 $k^{2\alpha}$  – thermal conductivity of the fractal material, [ $\text{Wm}^{-\alpha}\text{K}^{-1}$ ]  
 $\tilde{L}_\alpha\{f(x)\}$  – Yang-Laplace transform of  $f(x)$   
 $\tilde{L}_\alpha^{-1}[f_s^{L,\alpha}(s)]$  – inverse version of Yang-Laplace transform of  $f_s^{L,\alpha}(s)$

$t$  – time, [s]  
 $u(x,t)$  – temperature function, [K]  
 $x$  – space co-ordinate, [m]

### Greek symbols

$\alpha$  – fractal dimensional order, [–]  
 $\rho_\alpha$  – density, [ $\text{kgm}^{-3}$ ]

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